

Measurement of $\Xi^- \rightarrow \Lambda \pi^-$ and $\Omega^- \rightarrow \Lambda K^-$ Decay Parameters

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BEACH 2004

Illinois Institute of Technology, Chicago, USA

June 28–July 03, 2004

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Why We Measure Decay Parameters?

- A nonleptonic hyperon decay, $B_i \rightarrow B_f X$, is described by three decay parameters (B and X represent a baryon and a spin-zero meson respectively):

$$\alpha_H = \frac{2\text{Re}(A_{J-\frac{1}{2}}^* A_{J+\frac{1}{2}})}{|A_{J-\frac{1}{2}}|^2 + |A_{J+\frac{1}{2}}|^2}, \quad \beta_H = \frac{2\text{Im}(A_{J-\frac{1}{2}}^* A_{J+\frac{1}{2}})}{|A_{J-\frac{1}{2}}|^2 + |A_{J+\frac{1}{2}}|^2}, \quad \gamma_H = \frac{|A_{J-\frac{1}{2}}|^2 - |A_{J+\frac{1}{2}}|^2}{|A_{J-\frac{1}{2}}|^2 + |A_{J+\frac{1}{2}}|^2}.$$

- For $J = \frac{1}{2}$ ($J = \frac{3}{2}$), $A_{J-\frac{1}{2}}$ and $A_{J+\frac{1}{2}}$ are amplitudes of S -wave (P -wave) and P -wave (D -wave), and correspond to parity-violating (parity-conserving) and parity-conserving (parity-violating) states respectively;

- $\Xi \rightarrow \Lambda \pi$ decays: $\alpha_\Xi = \frac{2\text{Re}(S^* P)}{|S|^2 + |P|^2}, \quad \beta_\Xi = \frac{2\text{Im}(S^* P)}{|S|^2 + |P|^2}, \quad \gamma_\Xi = \frac{|S|^2 - |P|^2}{|S|^2 + |P|^2}.$

- $\Lambda \rightarrow p \pi$ decays: $\alpha_\Lambda = \frac{2\text{Re}(S^* P)}{|S|^2 + |P|^2},$

- $\Omega \rightarrow \Lambda K$ decays: $\alpha_\Omega = \frac{2\text{Re}(P^* D)}{|P|^2 + |D|^2}.$

$\Xi \rightarrow \Lambda\pi$ decays

- CP -violation in the $\Xi \rightarrow \Lambda\pi$ decay can be probed by the asymmetry

$$A_{\Xi} = - \underbrace{\tan(\delta_p - \delta_s)}_{\text{strong}} \underbrace{\sin(\phi_p - \phi_s)}_{\text{weak}}. \quad (1)$$

If CP -violating weak phases are negligible, $\delta_p - \delta_s$ can be determined through

$$\frac{\beta_{\Xi}}{\alpha_{\Xi}} = \tan(\delta_p - \delta_s) \quad (2)$$

Thus measuring β_{Ξ} and α_{Ξ} is equivalent to measuring $\delta_p - \delta_s$.

- **Theoretical Predictions**

Theory	$\delta_p - \delta_s$ (degree)
R. Nath (1965)	$\delta_p - \delta_s \approx 16$
M. Lu et al. (1994)	$\delta_p - \delta_s = -1.7$
A. Kamal (1998)	$-4.2 < \delta_p - \delta_s < -1.4$
A. Datta et al.(1998)	$-3.3 < \delta_p - \delta_s < 0.9$
U. Meissner and J. Oller (2001)	$-2.8 < \delta_p - \delta_s < -1.7$
J. Tandean et al. (2001)	$-3.9 < \delta_p - \delta_s < 7.8$

$\Omega \rightarrow \Lambda K$ decays

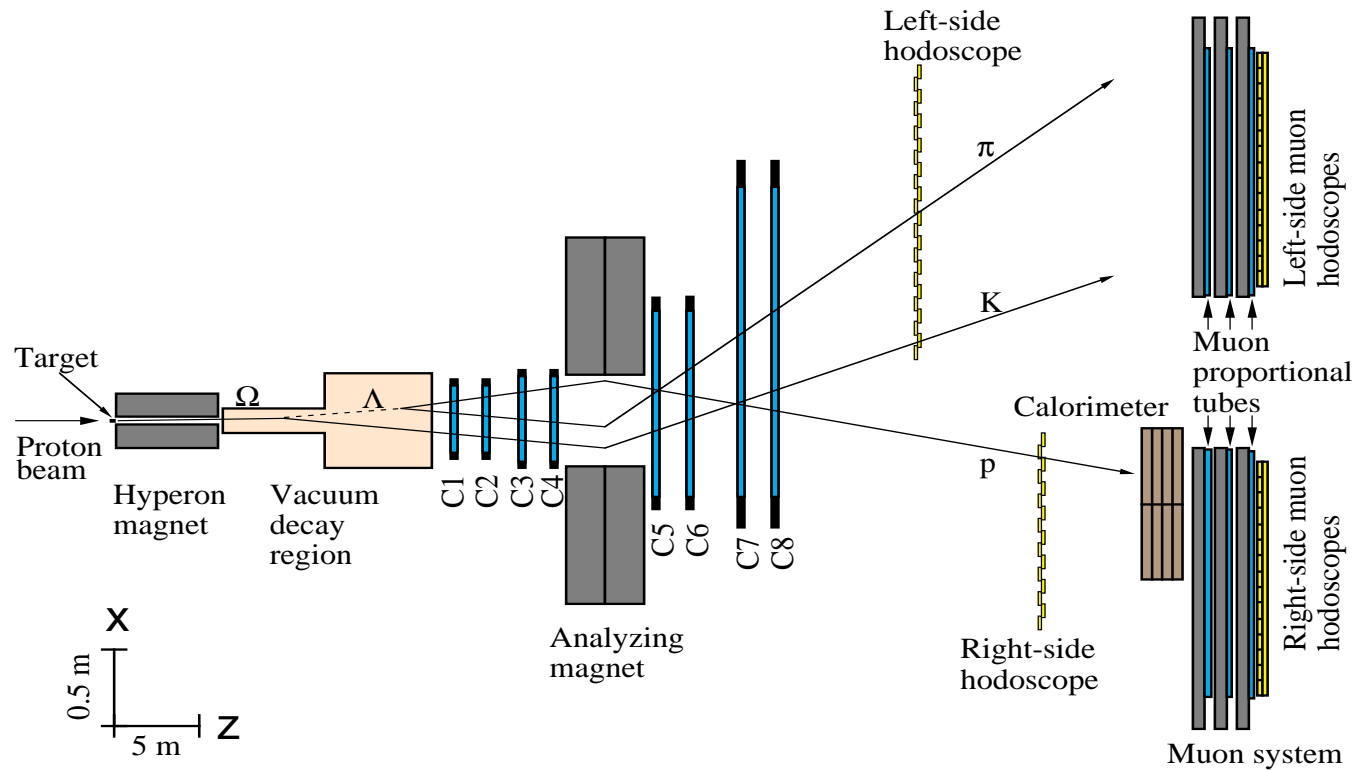
- Theoretically, D -wave is kinematically suppressed and α_Ω is expected to be zero (M. Suzuki, 1964). Non-zero α_Ω indicates parity violation in this decay.
- A difference between $|\alpha_\Omega|$ and $|\alpha_{\bar{\Omega}}|$ would be evidence of CP violation.
- Previous experimental results for α_Ω are **consistent with zero**.

Experiment	Year	Events	α_Ω
CERN-SPS	1984	12,000	-0.025 ± 0.028
FNAL E620	1988	1,743	-0.034 ± 0.079
FNAL E756	1998	6,953	-0.028 ± 0.047
PDG Average			-0.026 ± 0.023

- *HyperCP* has collected the largest samples of hyperon decays in the world which allows precise measurements of decay parameters.

The *HyperCP* (E871) Spectrometer

- Protons on target = (7 ~ 8) GHz
- Secondary beam intensity = (10 ~ 15) MHz
- Trigger rate = (50 ~ 80) KHz
- Charge of Sec. beam selected by the sign of Hyperon magnet



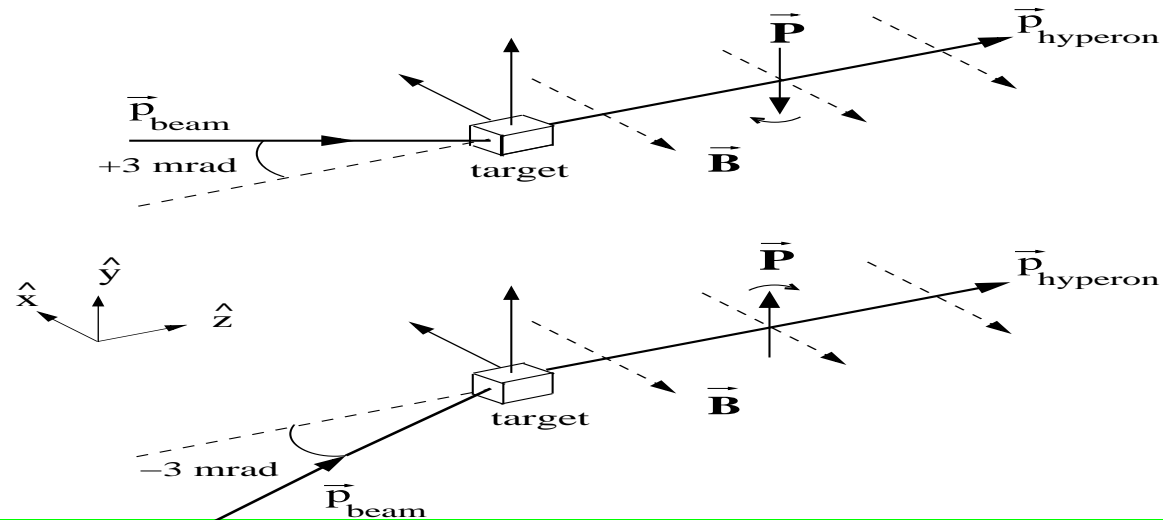
Measuring $\beta_{\Xi}, \gamma_{\Xi}$

- We use polarized $\Xi \rightarrow \Lambda\pi, \Lambda \rightarrow p\pi$ decays, where the angular distribution of the joint decay has the form

$$\frac{d^2N}{d\Omega_{\Lambda}d\Omega_p} = \frac{1}{(4\pi)^2} (1 + \alpha_{\Xi} \vec{P}_{\Xi} \cdot \hat{\Lambda}) (1 + \alpha_{\Lambda} \vec{P}_{\Lambda} \cdot \hat{p}), \quad (3)$$

where

$$\vec{P}_{\Lambda} = \frac{(\alpha_{\Xi} + \vec{P}_{\Xi} \cdot \hat{\Lambda}) \hat{\Lambda} + \beta_{\Xi} (\vec{P}_{\Xi} \times \hat{\Lambda}) + \gamma_{\Xi} \hat{\Lambda} \times (\vec{P}_{\Xi} \times \hat{\Lambda})}{1 + \alpha_{\Xi} \vec{P}_{\Xi} \cdot \hat{\Lambda}}. \quad (4)$$

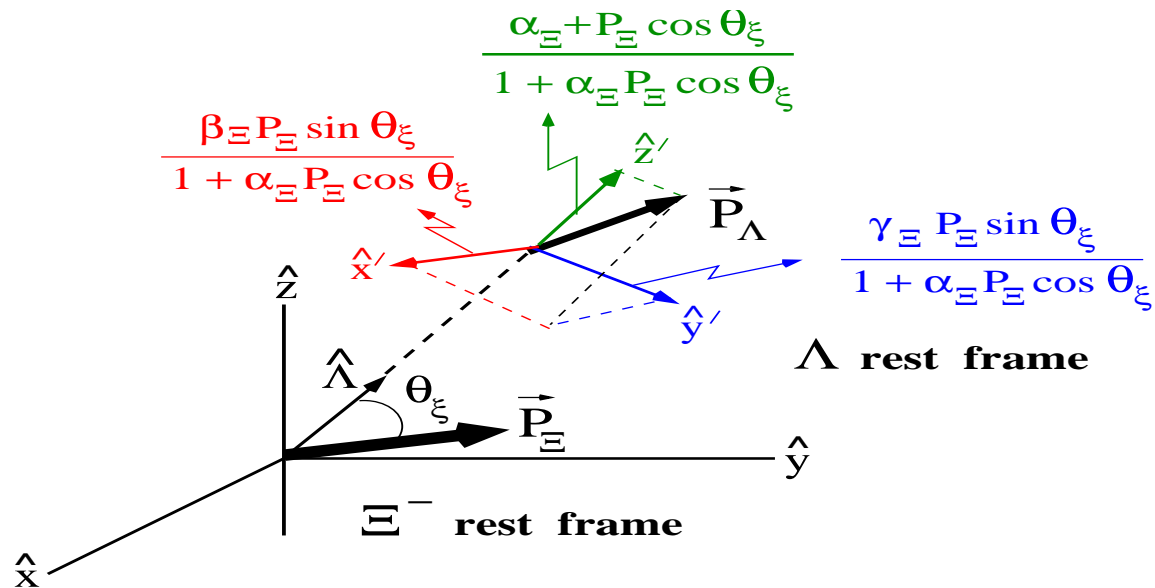


- In the coordinate system defined as

$$\hat{z}' = \hat{\Lambda}, \quad \hat{x}' = \frac{\vec{P}_{\Xi} \times \hat{\Lambda}}{|\vec{P}_{\Xi} \times \hat{\Lambda}|}, \quad \hat{y}' = \hat{z}' \times \hat{x}', \quad (5)$$

Λ polarization is determined as

$$\vec{P}_{\Lambda} = \left(\frac{\beta_{\Xi} P_{\Xi} \sin \theta_{\xi}}{1 + \alpha_{\Xi} P_{\Xi} \cos \theta_{\xi}}, \frac{\gamma_{\Xi} P_{\Xi} \sin \theta_{\xi}}{1 + \alpha_{\Xi} P_{\Xi} \cos \theta_{\xi}}, \frac{\alpha_{\Xi} + P_{\Xi} \cos \theta_{\xi}}{1 + \alpha_{\Xi} P_{\Xi} \cos \theta_{\xi}} \right)$$



- Proton angular distributions with respect to \hat{z}' , \hat{x}' , and \hat{y}' in Λ rest frame are respectively as

$$\frac{d^2 N}{d\Omega_\Lambda d \cos \theta_{pz'}} = \frac{1}{8\pi} [(1 + \alpha_\Xi \vec{P}_\Xi \cdot \hat{\Lambda}) + \alpha_\Lambda (\alpha_\Xi + \vec{P}_\Xi \cdot \hat{\Lambda}) \cos \theta_{pz'}], \quad (6)$$

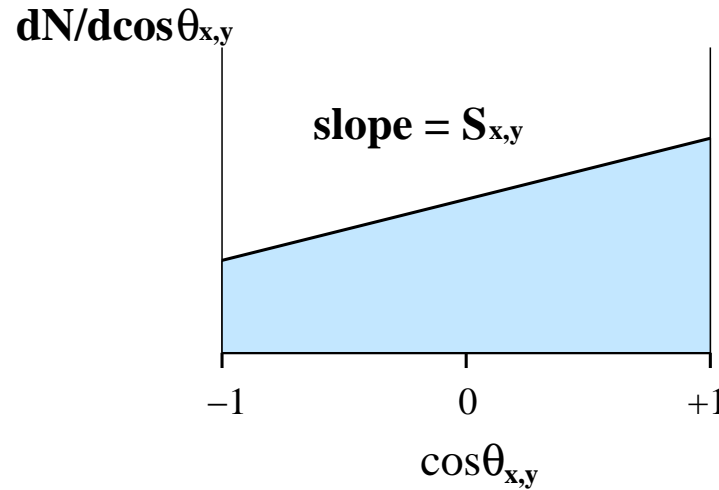
$$\frac{dN}{d \cos \theta_{px'}} = \frac{1}{2} \left(1 + \underbrace{\frac{\pi}{4} \alpha_\Lambda \beta_\Xi P_\Xi}_{S_x} \cos \theta_{px'} \right), \quad (7)$$

$$\frac{dN}{d \cos \theta_{py'}} = \frac{1}{2} \left(1 + \underbrace{\frac{\pi}{4} \alpha_\Lambda \gamma_\Xi P_\Xi}_{S_y} \cos \theta_{py'} \right), \quad (8)$$

where $\cos \theta_{pz'} = \hat{p} \cdot \hat{z}'$, $\cos \theta_{px'} = \hat{p} \cdot \hat{x}'$, $\cos \theta_{py'} = \hat{p} \cdot \hat{y}'$.

Eqs. (6), (7), and (8) allow us to measure \vec{P}_Ξ , β_Ξ , and γ_Ξ .

- β_{Ξ} and γ_{Ξ} can be extracted from the slopes of Eqs. (7) and (8), S_x and S_y .



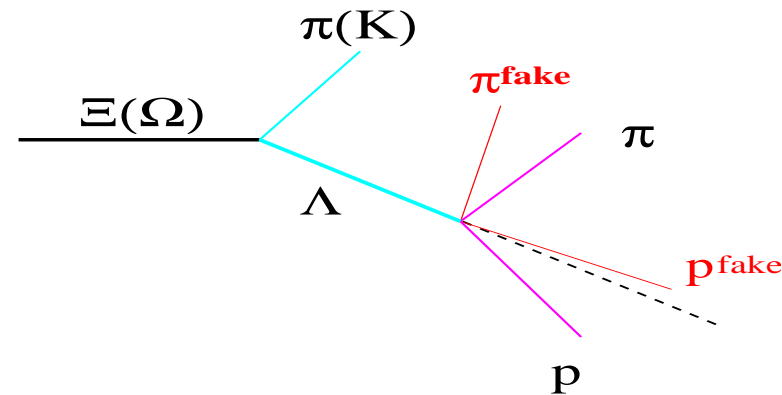
- Since $\alpha_{\Xi}^2 + \beta_{\Xi}^2 + \gamma_{\Xi}^2 = 1$, $(\phi_{\Xi}, \beta_{\Xi}, \gamma_{\Xi})$ is alternatively used to describe $\Xi \rightarrow \Lambda\pi$ decays

$$\beta_{\Xi} = \sqrt{1 - \alpha_{\Xi}^2} \sin \phi_{\Xi}, \quad \gamma_{\Xi} = \sqrt{1 - \alpha_{\Xi}^2} \cos \phi_{\Xi}. \quad (9)$$

where

$$\tan \phi_{\Xi} = \frac{\beta_{\Xi}}{\gamma_{\Xi}} = \frac{S_x}{S_y} \quad (10)$$

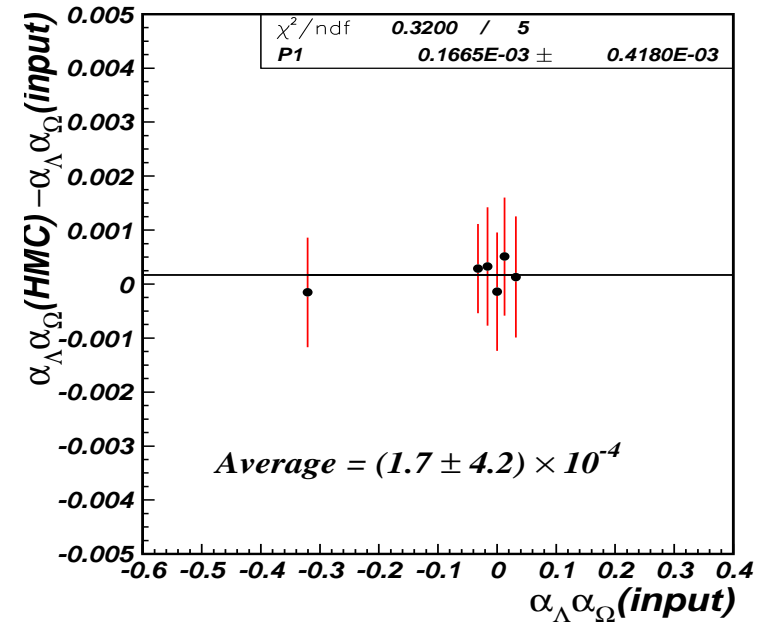
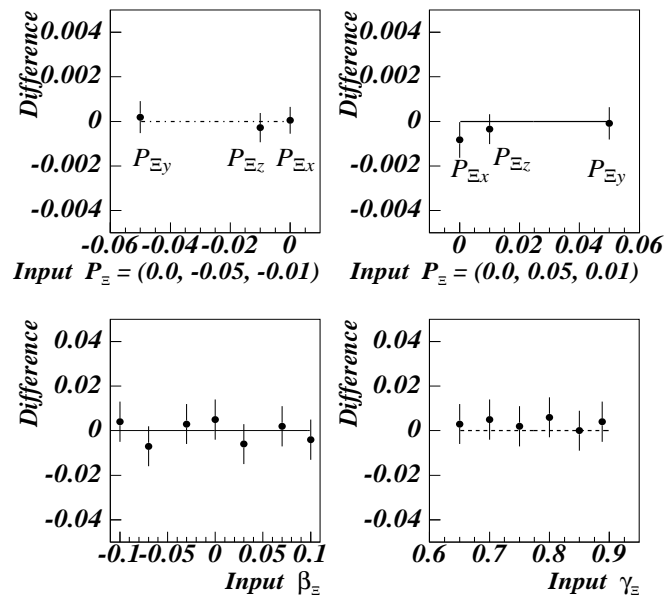
Hybrid Monte-Carlo Method (HMC)



- **In reality**, $\cos \theta$ distribution is distorted by the spectrometer acceptance, which can be measured by **Monte-Carlo**.
- Take all variables from each real event except $\cos \theta$.
- Generate **HMC events** with uniform $\cos \theta$.
- Let all the **HMC events** go through the software spectrometer, triggers, etc.; hence determine the acceptance.
- Weight **HMC events** to match the angular distributions of real and HMC events to obtain \vec{P}_Ξ , β_Ξ , γ_Ξ , and $\alpha_\Omega \alpha_\Lambda$.

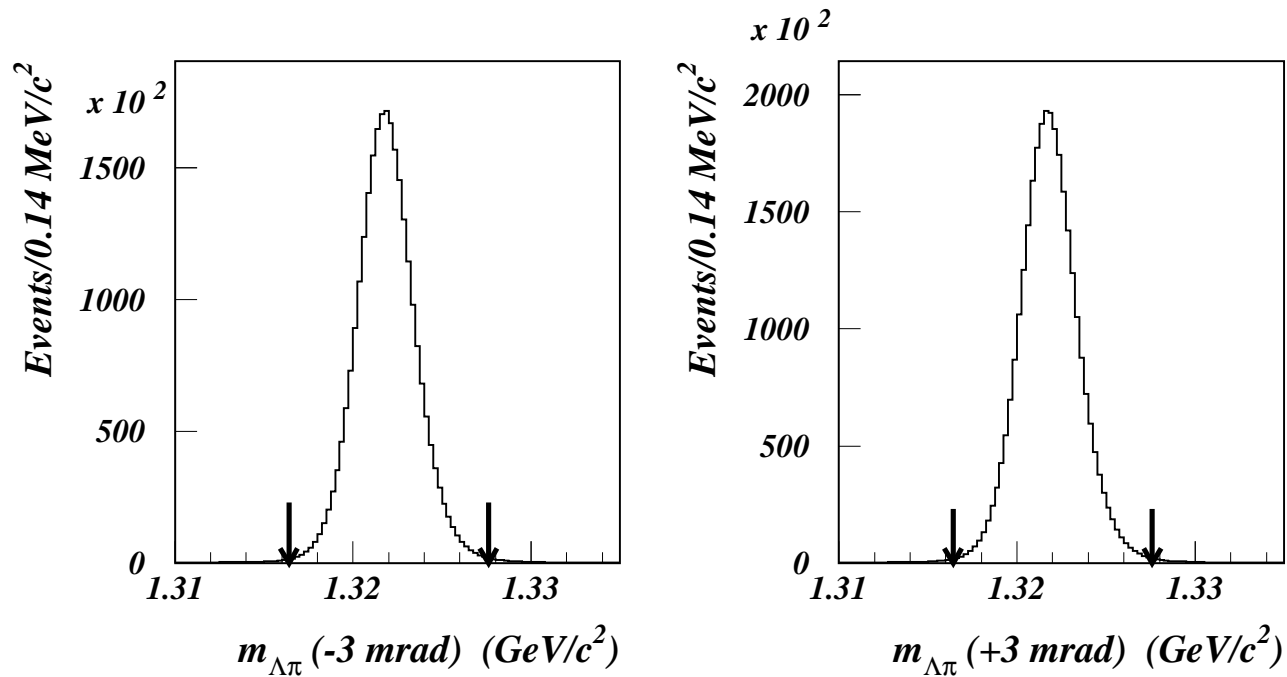
Validation of HMC

- Generated Monte-Carlo samples with different input values of \vec{P}_Ξ , β_Ξ , γ_Ξ , and $\alpha_\Omega\alpha_\Lambda$.
- Analyzed these Monte-Carlo samples with HMC to obtain the measurement values of \vec{P}_Ξ , β_Ξ , γ_Ξ , and $\alpha_\Omega\alpha_\Lambda$, and compare the difference between measurement and input values.

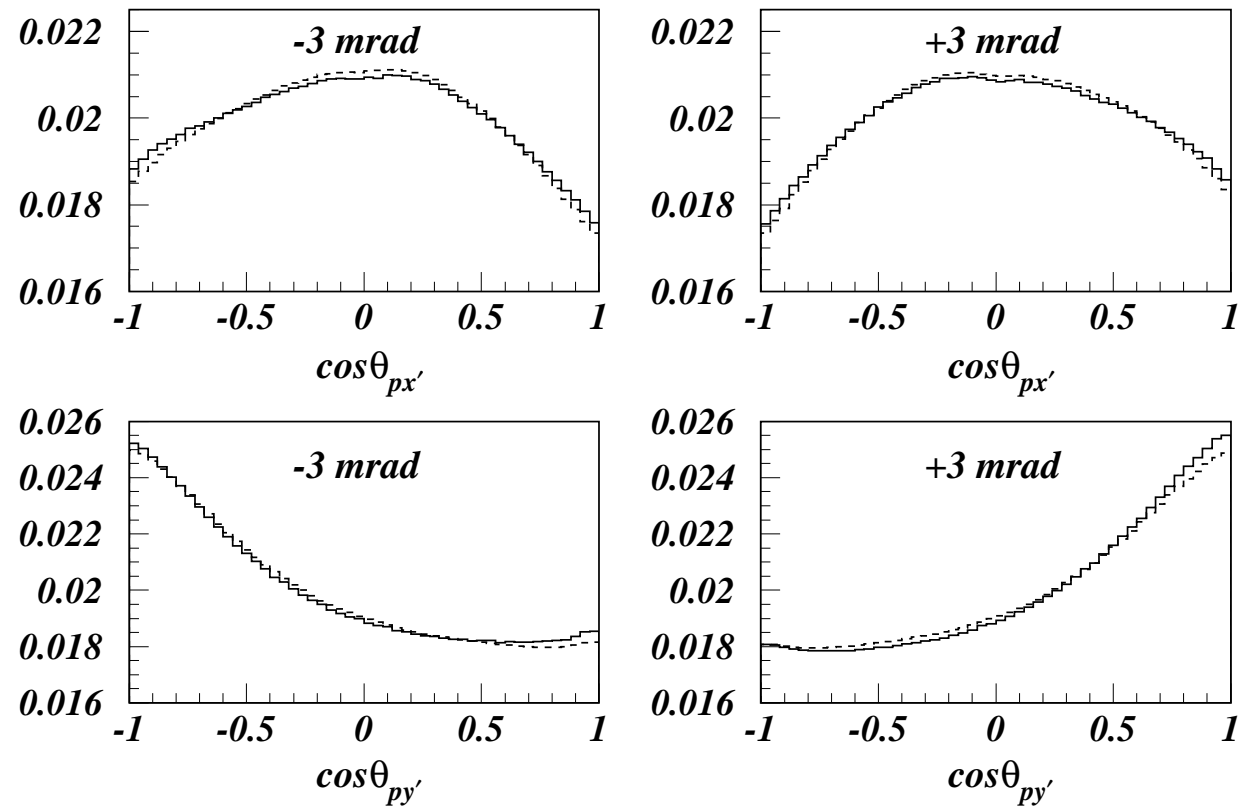


Results of ϕ_{Ξ} , β_{Ξ} and γ_{Ξ}

- 60 million -3 mrad and 72 million $+3 \text{ mrad}$ of polarized Ξ^{-} events after all event selection cuts were applied for the measurements.



- **Proton $\cos \theta_{px'}$ and $\cos \theta_{py'}$ distributions for data and HMC**



• **Results of ϕ_{Ξ} , β_{Ξ} and γ_{Ξ}**

p_{Ξ} (GeV/c)	S_x	S_y	ϕ_{Ξ} (degree)
139	-0.00037 ± 0.00047	0.01191 ± 0.00041	-1.77 ± 2.28
152	-0.00046 ± 0.00047	0.01447 ± 0.00038	-1.81 ± 1.88
162	-0.00038 ± 0.00041	0.01557 ± 0.00035	-1.39 ± 1.49
173	-0.00074 ± 0.00040	0.01880 ± 0.00036	-2.26 ± 1.22
191	-0.00123 ± 0.00040	0.02109 ± 0.00040	-3.33 ± 1.08
Average			-2.39 ± 0.64

Using $\alpha_{\Xi} = -0.458 \pm 0.012$ from PDG, and

$$\beta_{\Xi} = \sqrt{1 - \alpha_{\Xi}^2} \sin \phi_{\Xi} \text{ and } \gamma_{\Xi} = \sqrt{1 - \alpha_{\Xi}^2} \cos \phi_{\Xi},$$

we get

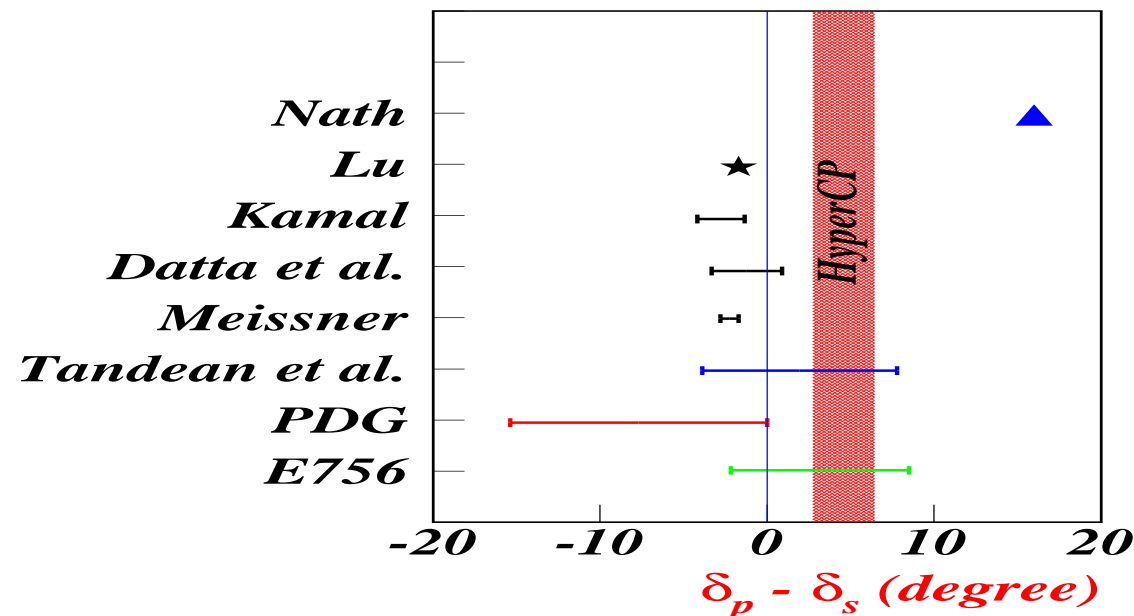
$$\beta_{\Xi} = -0.037 \pm 0.011(stat) \pm 0.010(syst),$$

$$\gamma_{\Xi} = 0.888 \pm 0.0004(stat) \pm 0.006(syst)$$

$$\delta_p - \delta_s = \tan^{-1}\left(\frac{\beta_{\Xi}}{\alpha_{\Xi}}\right) = [4.6 \pm 1.4(stat) \pm 1.2(syst)]^{\circ}$$

• Comparison With Other Results

- Chiral perturbation theory seems to have the wrong sign and small magnitude.
- $\delta_p - \delta_s$ is about a factor of 0.6 relative to the one of $p\pi$ scattering.
- CPV in $\Xi \rightarrow \Lambda\pi$ is smaller than that in $\Lambda \rightarrow p\pi$ but not negligible if strong phase-shift difference is considered only.

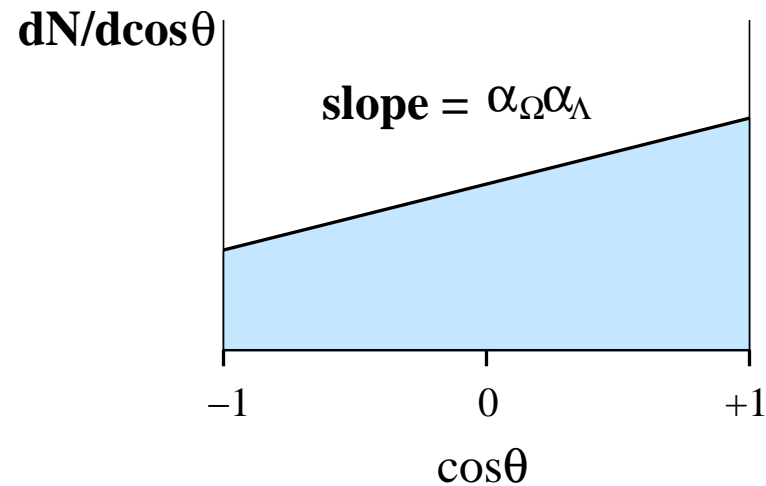


Measuring α_Ω

- For $\Omega \rightarrow \Lambda K$, $\Lambda \rightarrow p\pi$ decays, the proton angular distributions with respect to \hat{z}' , \hat{x}' , and \hat{y}' in Λ rest frame have the same form as those of the Ξ , except that those with respect to \hat{x}' and \hat{y}' have additional **tensor polarization** terms because $J = \frac{3}{2}$.
- To measure α_Ω , we use **unpolarized** $\Omega \rightarrow \Lambda K$ events. Hence the terms associated with β_Ω and γ_Ω vanish, and the proton angular distribution is simplified as ($\cos \theta \equiv \cos \theta_{pz'}$)

$$\frac{dN}{d \cos \theta} = \frac{1}{2} (1 + \underbrace{\alpha_\Omega \alpha_\Lambda}_{S_m} \cos \theta). \quad (11)$$

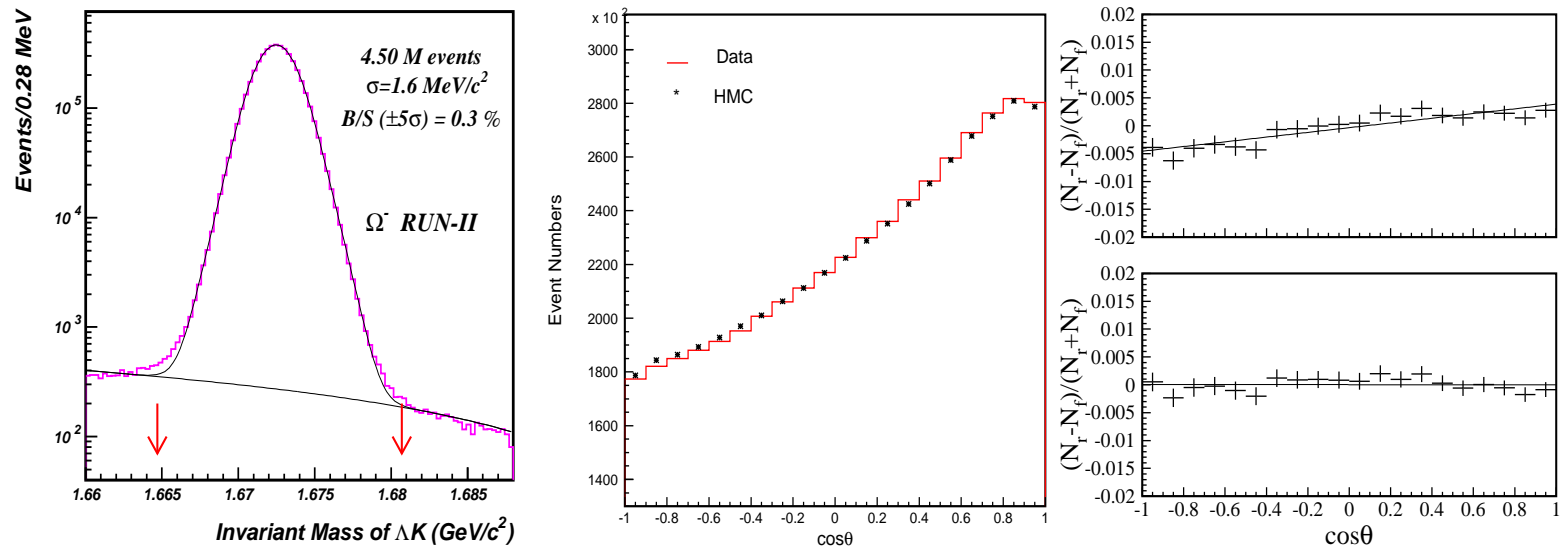
- Measuring the slope of the proton $\cos \theta$ distribution gives us $\alpha_\Omega \alpha_\Lambda$



- α_Ω can be extracted by using the parameter α_Λ of $\Lambda \rightarrow p\pi$ decays from PDG, $\alpha_\Lambda = 0.642 \pm 0.013$.

Results of α_Ω

- After all event selection cuts, 4.5 million from RUN-II $\Omega^- \rightarrow \Lambda K^-$ events were used to measure α_Ω .



- Slope of proton angular distribution S_m was extracted using HMC,

$$S_m = (1.16 \pm 0.12) \times 10^{-2}$$

- **Background slope S_b was measured from mass sidebands,**

$$S_b = (7.17 \pm 3.04) \times 10^{-2}.$$

- **Using $\alpha_\Omega \alpha_\Lambda = \frac{N_m}{N_s} S_m - \frac{N_b}{N_s} S_b$ to subtract background, we get**

$$\alpha_\Omega \alpha_\Lambda = [1.14 \pm 0.12(stat)] \times 10^{-2}.$$

- **Systematic Errors**

Systematic Study	$ \alpha_\Omega \alpha_\Lambda - (\alpha_\Omega \alpha_\Lambda)_{basic} $
π_Λ decays in flight	negligible
Event selection cuts	$0.09 \times 10^{-2} (0.73\sigma)$
Inefficiencies of spectrometer	negligible
BM109 field	negligible
HMC validation	$0.04 \times 10^{-2} (0.33\sigma)$
Total	0.10×10^{-2}

- **Results of $\alpha_\Omega\alpha_\Lambda$ and α_Ω**

Together with the results from the analysis of 1 million events on RUN-I data, our results are:

$$\alpha_\Omega\alpha_\Lambda = [1.18 \pm 0.29(\text{stat})] \times 10^{-2} \text{ (RUN-I),}$$

$$\alpha_\Omega\alpha_\Lambda = [1.14 \pm 0.12(\text{stat}) \pm 0.10(\text{syst})] \times 10^{-2} \text{ (RUN-II).}$$

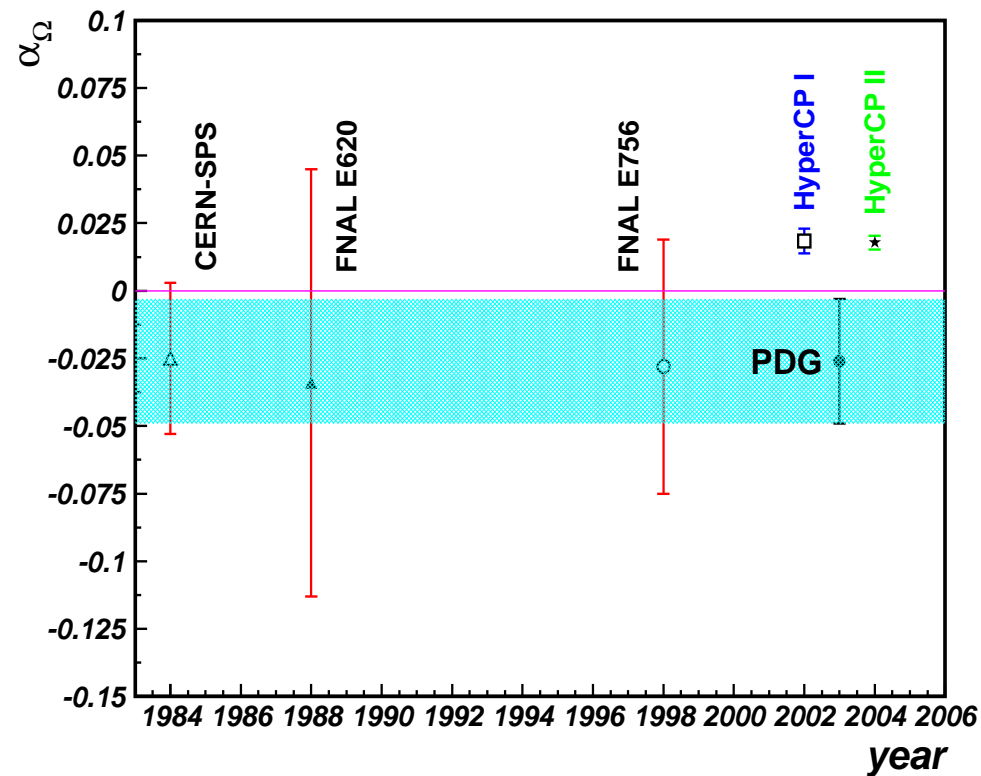
Using PDG value $\alpha_\Lambda = 0.642 \pm 0.013$, we get

$$\alpha_\Omega = [1.84 \pm 0.46(\text{stat})] \times 10^{-2} \text{ (RUN-I),}$$

$$\alpha_\Omega = [1.78 \pm 0.19(\text{stat}) \pm 0.16(\text{syst})] \times 10^{-2} \text{ (RUN-II).}$$

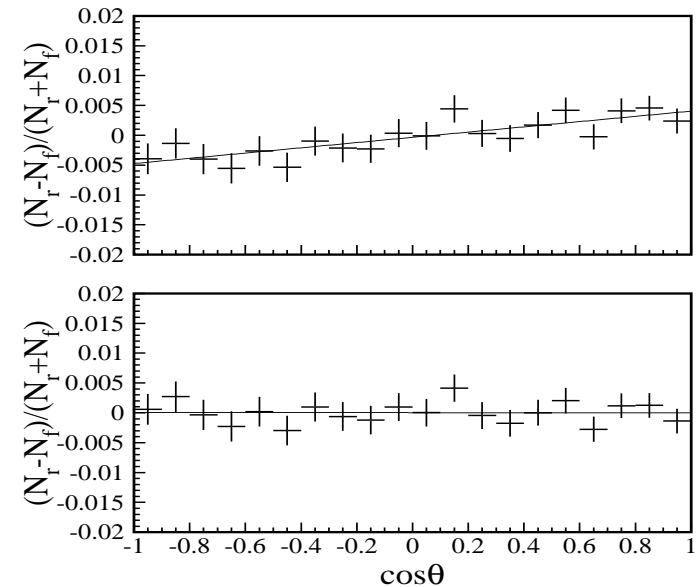
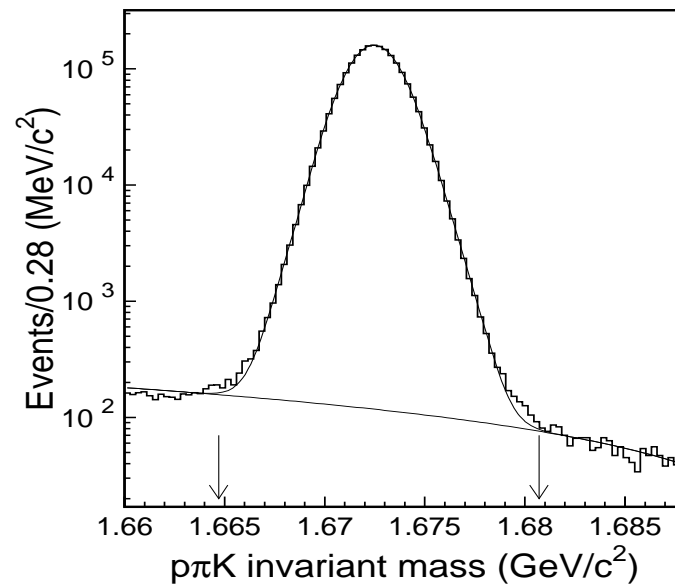
- **Comparing With Other Experimental Results**

- α_Ω from *HyperCP* has a positive sign, which is different from PDG value and other experiments.
- The result of α_Ω from *HyperCP* is nonzero (7.2σ).



Preliminary Measurement of $\alpha_{\bar{\Omega}}$ and CPV Test

- Using the same code and event selection cuts, we have analyzed 1.9 million $\bar{\Omega}^+ \rightarrow \bar{\Lambda} K^+, \bar{\Lambda} \rightarrow \bar{p} \pi^+$ events.



- **Preliminary Results of $\alpha_{\bar{\Omega}}\alpha_{\bar{\Lambda}}$ and $\alpha_{\bar{\Omega}}$**

$$\alpha_{\bar{\Omega}}\alpha_{\bar{\Lambda}} = [1.16 \pm 0.18(\text{stat})] \times 10^{-2}.$$

Using the PDG value of $\alpha_{\bar{\Lambda}} = -0.642 \pm 0.013$, $\alpha_{\bar{\Omega}}$ was extracted:

$$\alpha_{\bar{\Omega}} = [-1.81 \pm 0.28(\text{stat})] \times 10^{-2}.$$

- **CPV Test For $\Omega \rightarrow \Lambda K$ Decays**

Using the measured values of $\alpha_{\Omega}\alpha_{\Lambda}$ and $\alpha_{\bar{\Omega}}\alpha_{\bar{\Lambda}}$, asymmetries are determined to be

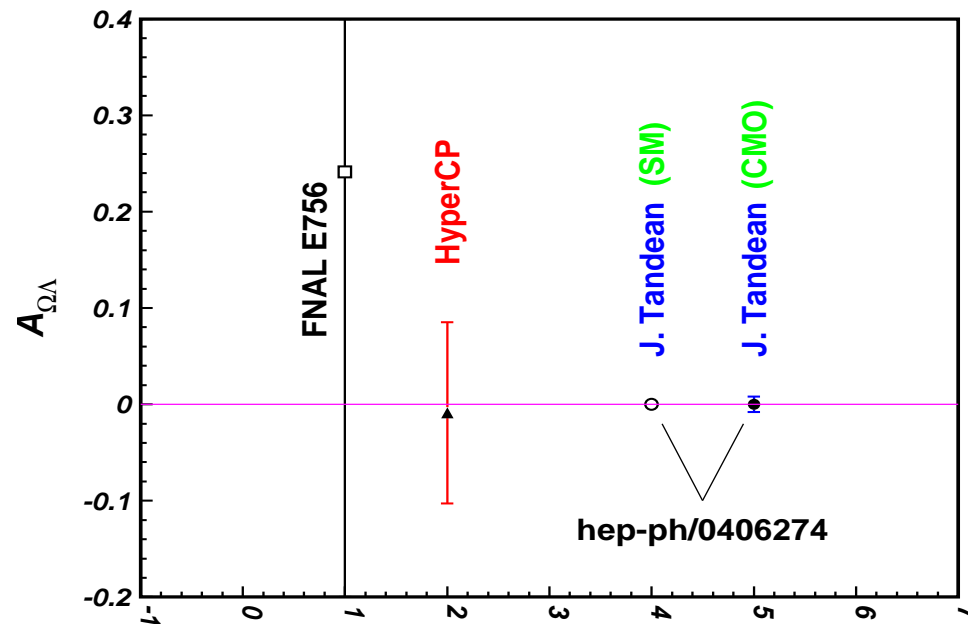
$$\delta_{\Omega\Lambda} \equiv \alpha_{\Omega}\alpha_{\Lambda} - \alpha_{\bar{\Omega}}\alpha_{\bar{\Lambda}} = [-0.02 \pm 0.22(\text{stat})] \times 10^{-2},$$

$$A_{\Omega\Lambda} \equiv \frac{\alpha_{\Omega}\alpha_{\Lambda} - \alpha_{\bar{\Omega}}\alpha_{\bar{\Lambda}}}{\alpha_{\Omega}\alpha_{\Lambda} + \alpha_{\bar{\Omega}}\alpha_{\bar{\Lambda}}} = [-0.87 \pm 9.41(\text{stat})] \times 10^{-2}.$$

CP is conserved within statistics

• $A_{\Omega\Lambda}$ From Other Measurements and Theoretical Predictions

- $A_{\Omega\Lambda}$ for **FNAL E756** in the plot is extrapolated from their measurements, $\alpha_{\Omega} = -0.028 \pm 0.047$ and $\alpha_{\overline{\Omega}} = +0.017 \pm 0.077$.
- In Tandean's predictions, **SM** stands for the standard model, and **CMO** stands for chromomagnetic-penguin operators which includes possible physics beyond the standard model.

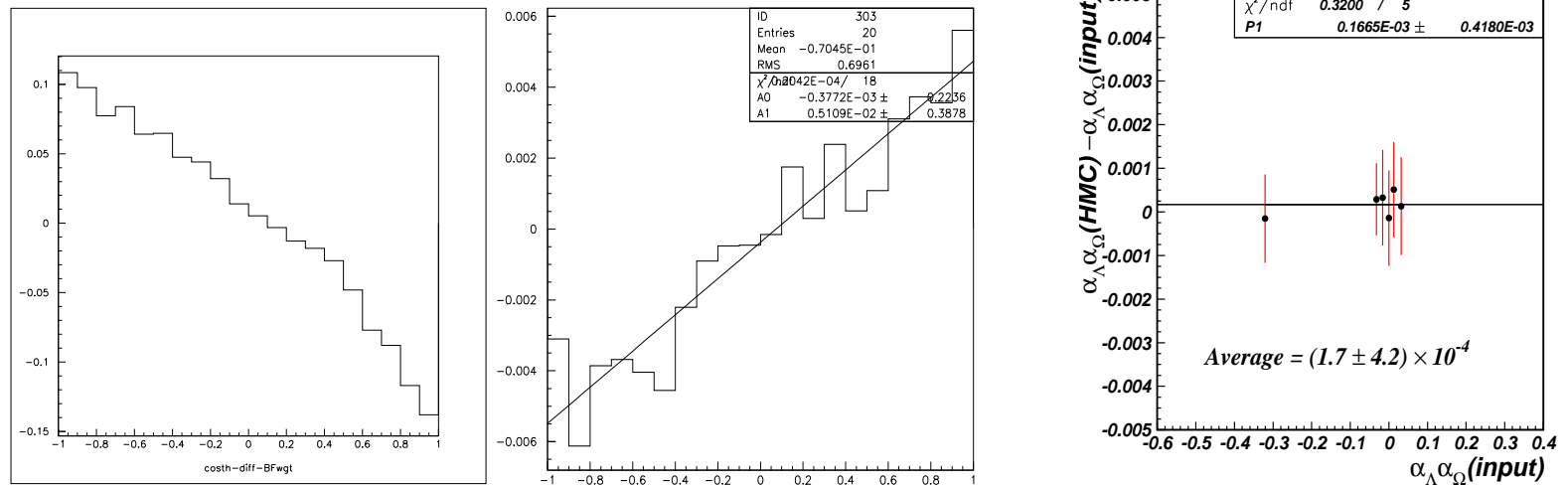


Conclusions

- With the largest sample of hyperon decays ever recorded, HyperCP has precisely measured the decay parameters: β_{Ξ} , γ_{Ξ} and α_{Ω} .
- The $\Lambda\pi$ scattering strong phase-shift difference deduced from the measured β_{Ξ} and PDG α_{Ξ} is nonzero and is comparable to the one of $p\pi$ scattering, which indicates CPV in $\Xi \rightarrow \Lambda\pi$ decays may not be suppressed compared to that in $\Lambda \rightarrow p\pi$ decays.
- We find the first evidence of a nonzero α_{Ω} , and hence parity violation in $\Omega \rightarrow \Lambda K$ decays.
- No CP -violation is observed in the joint decays of $\Omega \rightarrow \Lambda K \rightarrow p\pi K$ at the level of 0.22×10^{-2} for $\delta_{\Omega\Lambda}$ and 9.41×10^{-2} for $A_{\Omega\Lambda}$.

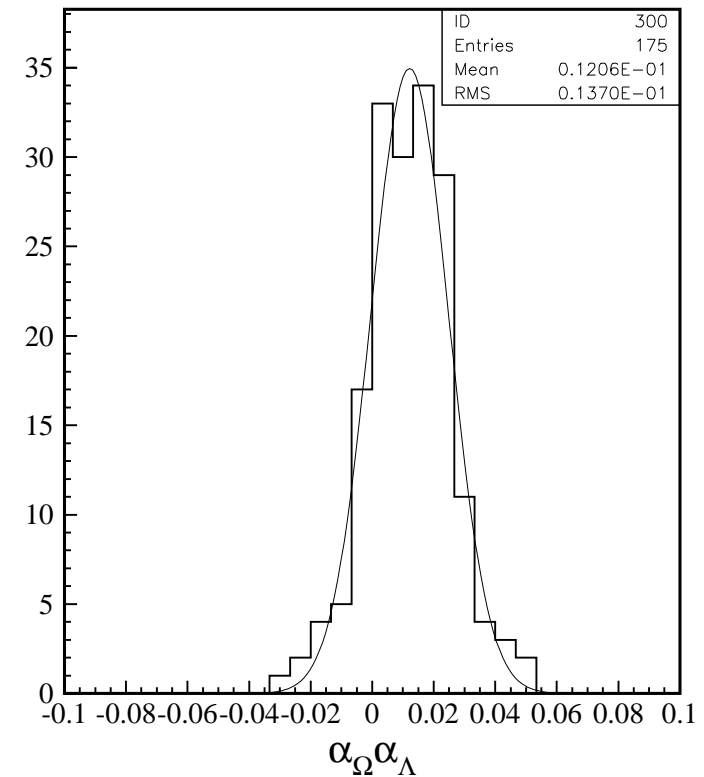
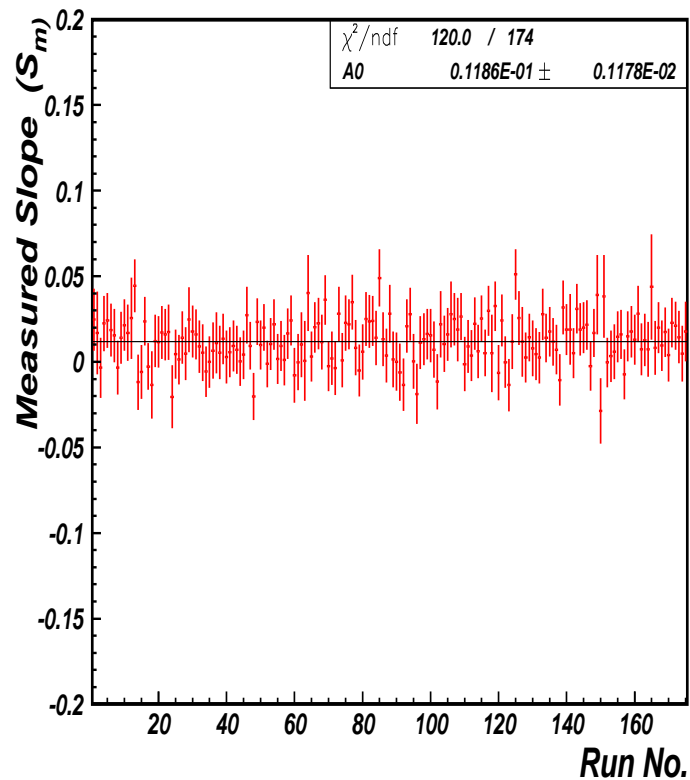
The Sign of α_Ω

- The same analysis code was used to analyze 78,000 events of the $\Xi^- \rightarrow \Lambda \pi^-$ decay, $\alpha_\Xi \alpha_\Lambda = -0.286 \pm 0.007$, both the sign and value are consistent with PDG.



- Monte-Carlo Study: HMC results consistent with input values both on the sign and value.

- Individual run by run analysis shows only 29 out of 175 runs have negative $\alpha_\Omega \alpha_\Lambda$ values but all with very large statistical errors compared with their values.



Systematic Study of β and γ

- Systematic Errors

Systematic Study	S_x ($\times 0.00019$)	S_y ($\times 0.00017$)
Accidental hits	0.28	0.06
Background	0.06	0.22
Detector efficiency	0.39	0.50
Precession angle	0.58	0.17
Event selection	0.61	0.64
Total error	0.97	0.86